

MATH 241
LECTURE 9

: More problems.

- Solve the following problem using separation of variables:

$$\left. \begin{aligned} u_{xx} + u_{yy} &= 0, \quad 0 < x < \infty, \quad 0 < y < H, \\ u_y(x, 0) &= 0 \\ u_y(x, H) &= 0 \\ u_x(0, y) &= f(y) \end{aligned} \right\}$$

Questions: 1) Show that there exists a solution if and only if $\int_0^H f(y) dy = 0$.

2) Is the solution uniquely defined?

Sol:

$$\text{Let } u(x, y) = F(x)G(y) \rightarrow \frac{-F''(x)}{F(x)} = \frac{G''(y)}{G(y)} = -\lambda.$$

BC:

$$u_y(x, 0) = 0 \Rightarrow G'(0) = 0,$$

$$u_y(x, H) = 0 \Rightarrow G'(H) = 0,$$

$$\text{ODE: } \left\{ \begin{aligned} G''(y) + \lambda G(y) &= 0 \\ G'(0) = 0 = G'(H) \end{aligned} \right. \rightarrow \left\{ \begin{aligned} \lambda &= \left(\frac{n\pi}{H}\right)^2, \quad n = 0, 1, \dots \\ G(y) &= \cos\left(\frac{n\pi}{H}y\right), \quad n = 0, 1, \dots \end{aligned} \right.$$

$$\textcircled{2} \quad F''(x) - \left(\frac{n\pi}{H}\right)^2 F(x) = 0$$

$$n=0 \Rightarrow F(x) = c_1 x + c_2$$

$$n \neq 0 \Rightarrow F(x) = c_3 e^{\frac{n\pi x}{H}} + c_4 e^{-\frac{n\pi x}{H}}$$

Remark: Physically valid solutions should be bounded as $x \rightarrow +\infty$. Therefore, we take $c_1 = c_3 = 0$.

$$\text{So, } F(x) = e^{-\frac{n\pi x}{H}}, \quad n=0, 1, \dots$$

$$\bullet \text{ Product solution: } u(x, y) = A \cos\left(\frac{n\pi y}{H}\right) e^{-\frac{n\pi x}{H}}$$

$\bullet$ Superposition principle:

$$u(x, y) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi y}{H}\right) e^{-\frac{n\pi x}{H}}.$$

$$\bullet A_n? \rightarrow u_x(0, y) = f(y)$$

$$u_x(x, y) = \sum_{n=0}^{\infty} A_n \cos\left(\frac{n\pi y}{H}\right) \left(-\frac{n\pi}{H}\right) e^{-\frac{n\pi x}{H}},$$

$$u_x(0, y) = \sum_{n=0}^{\infty} A_n \frac{-n\pi}{H} \cos\left(\frac{n\pi y}{H}\right) = f(y).$$

Using the orthogonality of cosines, we find that

$$\int_0^H f(y) \cos\left(\frac{n\pi y}{H}\right) dy = A_n \frac{-n\pi}{H} \frac{H}{2}, \quad n=1,2,\dots \Rightarrow$$

$$\Rightarrow A_n = \frac{-2}{\pi} \int_0^H f(y) \cos\left(\frac{n\pi y}{H}\right) dy, \quad n=1,2,\dots$$

Notice what happens when $n=0$:

$$\int_0^H f(y) dy = \sum_{n=0}^{\infty} A_n \frac{-n\pi}{H} \int_0^H \cos\left(\frac{n\pi y}{H}\right) dy = \sum_{n=1}^{\infty} A_n \frac{-n\pi}{H} \int_0^H \cos\left(\frac{n\pi y}{H}\right) dy$$

$= 0$

(A dashed arrow points from the $n=0$ term to the $n=1$ term in the sum.)

so we obtain the condition we were asked to show.

Therefore, the solution is

$$u(x,y) = A_0 + \sum_{n=1}^{\infty} A_n \cos\left(\frac{n\pi y}{H}\right) e^{-\frac{n\pi x}{H}}, \quad \text{with}$$

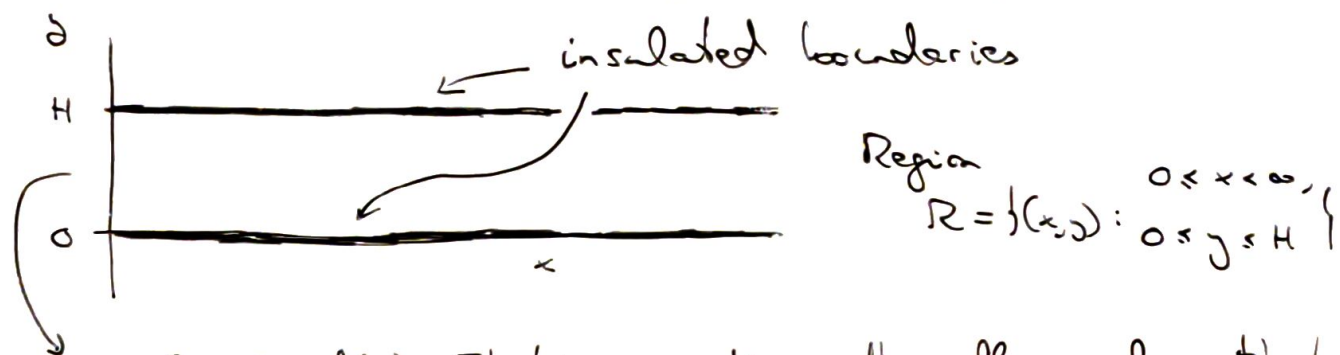
$$A_n = \frac{-2}{\pi} \int_0^H f(y) \cos\left(\frac{n\pi y}{H}\right) dy \quad \text{for } n=1,2,\dots$$

Thus, we see that the solution is not uniquely defined.

The constant A_0 is arbitrary.

→ What is happening?

Let's think about the physics of the problem:



→ Remember that the heat equation over a 2-d region is $u_t = k \Delta u = k(u_{xx} + u_{yy})$, and therefore we can think of the Laplace equation $\Delta u = 0$ as the PDE that describes the equilibrium temperature over that region.

↳ For an equilibrium to exist, we need that the "net flux" is zero: otherwise, the region would receive (or give) energy, and the temperature would increase (or decrease). Mathematically, the condition

$$\int_0^H f(y) dy = 0 \quad \text{precisely states that the net (or average)}$$

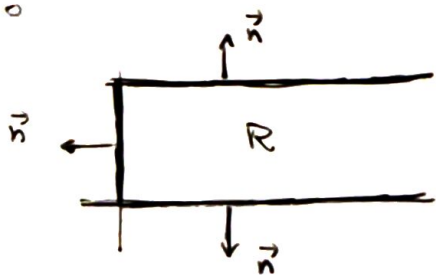
flux is zero along the boundary $x=0, 0 \leq y \leq H$.

Since the B.C., with the condition $\int_0^H f(y) dy = 0$, state that there is no _{net} energy going in or out, any constant temperature is an equilibrium solution (notice that there isn't any Dirichlet B.C.).

→ If we were given an initial condition, then we'd be able to determine the solution, by conservation of energy. Indeed, let's see it mathematically. Assume that $u(x, y, 0) = g(x, y)$, with $\iint_R g(x, y) dx dy = 0$. Consider the solution to $u_t = u_{xx} + u_{yy}$ with the same B.C.

Then,

$$\begin{aligned} \frac{d}{dt} \iint_R u(x, y, t) dx dy &= \iint_R \Delta u(x, y, t) dx dy = \iint_R \nabla \cdot \nabla u(x, y, t) dx dy = \\ &= \int_{\partial R} \vec{n} \cdot \nabla u(x, y, t) dx dy = \int_0^H (-u_x(0, y)) dy + \int_0^a u_y(x, H) dx + \\ &+ \int_0^a (-u_y(x, 0)) dx = 0 \quad (\text{because } \int_0^H u_x(0, y) dy = \int_0^H f(y) dy = 0). \end{aligned}$$



Therefore,
$$\iint_R u(x, y, t) dx dy = \iint_R u(x, y, 0) dx dy = \iint_R g(x, y) dx dy.$$

This holds for all $t \geq 0$, so in particular at equilibrium:

$$\iint_R u(x, y) dx dy = \iint_R g(x, y) dx dy = 0,$$

The solution we found
for the Laplace equation

we can compute the first integral:

$$\begin{aligned} \iint_R u(x, y) dx dy &= \iint_R A_0 dx dy + \sum_{n=1}^{\infty} A_n \int_0^H \cos\left(\frac{n\pi y}{H}\right) dy \int_0^a e^{\frac{-n\pi x}{H}} dx = \\ &= \iint_R A_0 dx dy = 0 \rightarrow A_0 = 0. \end{aligned}$$

[Problem 3, pdf] Consider $\begin{cases} x^2 \phi''(x) + x \phi'(x) + \lambda \phi(x) = 0, & 1 < x < e, \\ \phi(1) = \phi(e) = 0. \end{cases}$

a) Find all the eigenvalues and eigenfunctions.

b) Construct a general solution to

$$x^2 u_{xx}(x, t) + x u_x(x, t) = u_t(x, t), \quad 1 < x < e, t > 0,$$

given that $u(1, t) = 0 = u(e, t)$ for $t \geq 0$.

Sol:

• $\lambda = \alpha^2 > 0$

Try with x^p : $x^2 p(p-1)x^{p-2} + x p x^{p-1} + \lambda x^p = 0 \Rightarrow$

$$\Rightarrow p^2 + \alpha^2 = 0 \Rightarrow p = \pm i\alpha$$

Thus, $\phi(x) = c_1 x^{i\alpha} + c_2 x^{-i\alpha}$ ($c_1, c_2 \in \mathbb{C}$).

Notice that $x^{i\alpha} = e^{\ln x^{i\alpha}} = e^{i\alpha \ln x} = \cos(\alpha \ln x) + i \sin(\alpha \ln x)$.

Doing the same with $x^{-i\alpha}$, we find that

$$\phi(x) = (a_1 + ib_1)(\cos(\alpha \ln x) + i \sin(\alpha \ln x)) + (a_2 + ib_2)(\cos(\alpha \ln x) - i \sin(\alpha \ln x)), \text{ where}$$

we are writing $c_1 = a_1 + ib_1$, $c_2 = a_2 + ib_2$, $a_1, a_2, b_1, b_2 \in \mathbb{R}$.

Since we are only interested in real solutions, we get

$$\phi(x) = (a_1 + a_2) \cos(|\lambda| \ln x) + (b_2 - b_1) \sin(|\lambda| \ln x).$$

Renaming the constant,

$$\phi(x) = c_3 \cos(|\lambda| \ln x) + c_4 \sin(|\lambda| \ln x).$$

[Remark: We don't need to show all these steps everytime.]

B.C.: $\phi(1) = 0 = c_3$

$$\phi(e) = c_4 \sin(|\lambda| \ln e) = 0 \Rightarrow |\lambda| = n\pi, n=1, 2, \dots$$

Thus, $\lambda = (n\pi)^2$, $n=1, 2, \dots$ are eigenvalues,

with $\phi(x) = c \cdot \sin(n\pi \ln x)$ eigenfunctions.

• $\lambda = 0$: $\frac{d}{dx}(x\phi'(x)) = 0 \Rightarrow x\phi'(x) = c, \Rightarrow$

$$\Rightarrow \phi(x) = c_1 \ln x + c_2$$

B.C.: $\phi(1) = c_2 = 0$
 $\phi(e) = c_1 = 0 \quad \left\{ \Rightarrow \lambda = 0 \text{ is not an eigenvalue.} \right.$

• $\lambda = -\alpha^2 < 0$: Try with $x^p \rightarrow p = \pm |\alpha|$.

$$\phi(x) = c_1 x^{|\alpha|} + c_2 x^{-|\alpha|}.$$

$$\begin{aligned} \text{B.C. } \phi(1) &= c_1 + c_2 = 0 \\ \phi(e) &= c_1 e^{|\alpha|} + c_2 e^{-|\alpha|} = 0 \end{aligned} \quad \left. \vphantom{\begin{aligned} \phi(1) &= c_1 + c_2 = 0 \\ \phi(e) &= c_1 e^{|\alpha|} + c_2 e^{-|\alpha|} = 0 \end{aligned}} \right\} \Rightarrow c_1 = c_2 = 0.$$

Therefore, there are not negative eigenvalues.

b) $x^2 u_{xx}(x, t) + x u_x(x, t) = u_t(x, t), \quad 1 < x < e, \quad t > 0,$
 $u(1, t) = 0 = u(e, t), \quad t \geq 0.$

Sol $u(x, t) = h(t) \phi(x) \Rightarrow$

$$\Rightarrow x^2 h(t) \phi''(x) + x h(t) \phi'(x) = h'(t) \phi(x) \Rightarrow$$

$$\Rightarrow x^2 \frac{\phi''(x)}{\phi(x)} + x \frac{\phi'(x)}{\phi(x)} = \frac{h'(t)}{h(t)} = -\lambda \Rightarrow$$

$$\Rightarrow \textcircled{1} h'(t) = -h(t)$$

$$\textcircled{2} x^2 \phi''(x) + x \phi'(x) + \lambda \phi(x) = 0$$

B.C. $u(1, t) = 0 \Rightarrow \phi(1) = 0$

$$u(e, t) = 0 \Rightarrow \phi(e) = 0.$$

Thus, from part a),

$$\lambda = (n\pi)^2, n=1,2,\dots \text{ with } \phi(x) = c \sin(n\pi h x).$$

- ④ $h(t) = e^{-t}.$

- Product solution: $u(x,t) = c \sin(n\pi h x) e^{-t}.$

- Superposition principle:

$$u(x,t) = \sum_{n=1}^{\infty} B_n \sin(n\pi h x) e^{-t}.$$

The B_n are undetermined since we don't have initial condition.

• Problem 1: Let $u(x, t)$ be the temperature in a rod, that satisfies

$$\begin{cases} u_t = 3u_{xx} + x^3, & 0 < x < 4, t > 0, \\ u_x(0, t) = 1 \\ u_x(4, t) = 5 \\ u(x, 0) = x. \end{cases}$$

The total energy is defined by $E(t) = \int_0^4 u(x, t) dx$.

Find $E(t)$.

Sol: Integrate the equation from $x=0$ to $x=4$:

$$\begin{aligned} \int_0^4 u_t(x, t) dx &= 3 \int_0^4 u_{xx}(x, t) dx + \int_0^4 x^3 dx \Rightarrow \\ \Rightarrow \frac{d}{dt} \int_0^4 u(x, t) dx &= 3 u_x(x, t) \Big|_0^4 + \frac{x^4}{4} \Big|_0^4 \Rightarrow \end{aligned}$$

$$\Rightarrow \frac{dE}{dt}(t) = 3(5-1) + 4^3 = 12 + 56 = 68 \Rightarrow$$

$$\Rightarrow E(t) = 68t + c.$$

Using the initial condition, we find c :

$$E(0) = c = \int_0^4 u(x,0) dx = \int_0^4 x dx = \left. \frac{x^2}{2} \right|_0^4 = 8.$$

Therefore,

$$E(t) = 68t + 8.$$

• Problem 4: Find $u(x,t)$: $u_t - u = 7x$ with $u(x,0) = 0$.

Sol: Notice that e^{-t} is an integrating factor:

$$\frac{\partial}{\partial t} (e^{-t} u(x,t)) = 7x e^{-t} \Rightarrow e^{-t} u(x,t) = -7x e^{-t} + c(x)$$

we are integrating in t !

$$\left(\text{Check: } \frac{\partial}{\partial t} [e^{-t} u = -7x e^{-t} + c(x)] \rightsquigarrow \frac{\partial}{\partial t} (e^{-t} u) = 7x e^{-t} \right)$$

Thus,

$$u(x,t) = -7x + c(x) e^t.$$

$$\text{Then, } u(x,0) = -7x + c(x) = 0 \Leftrightarrow c(x) = 7x.$$

$$\text{In conclusion, } u(x,t) = 7x(e^t - 1).$$