

Poisson's Equation

The Poisson equation is the nonhomogeneous version of the Laplace equation, which can be interpreted as the time-independent version of the heat equation.

- We will see how to use the method of eigenfunction expansion. Since there is no time variable, we won't obtain an ODE as in lecture 19.

Problem: PDE: $\Delta u = Q$ in Ω } Ω some domain
 BC: $u = g$ on $\partial\Omega$ } in $\mathbb{R}^2$ (or $\mathbb{R}^3$).

[Q, g are given functions]

we could consider different BC in different parts of the boundary.

Solution:

- First, consider the associated modified homogeneous

problem: $\Delta \phi = -\lambda \phi$ in Ω (Helmholtz equation)
 $\phi = 0$ on $\partial\Omega$

same type of BC as in the original problem.

This eigenvalue problem will have solutions λ_i, ϕ_i .

Remark: For Ω a rectangle, circle, cylinder, sphere (or modified versions of them) we have already solved the Helmholtz equation with given BCs, having explicit λ_i, ϕ_i .

- Once we have λ_i, ϕ_i , we let our solution be an expansion in these eigenfunctions:

$$u = \sum_i c_i \phi_i, \text{ for some constants } c_i \text{ to be determined.}$$

▮

Remark: We might be tempted to say here: "Since

$u = \sum_i c_i \phi_i$, we substitute into the PDE to find

that $\Delta u = \cancel{\sum_i c_i \Delta \phi_i} = Q \rightarrow - \sum_i c_i \lambda_i \phi_i = Q \rightarrow$

$\xrightarrow{\text{orthogonality}} -c_n \lambda_n \int_{\Omega} (\phi_n)^2 = \int_{\Omega} \phi_n Q \rightarrow c_n = - \frac{1}{\lambda_n} \frac{\int_{\Omega} \phi_n Q}{\int_{\Omega} (\phi_n)^2}.$

→ The idea is "almost" right, but notice that we haven't used the BCs at all! Something must be wrong...

→ We cannot differentiate term by term on $u = \sum_i c_i \phi_i$, because u and ϕ_i do not satisfy the same BCs.

[In lect. 19 we assumed the BCs for u were homogeneous, because in lect. 18 we showed that we could always make the BCs homogeneous by changing the force].



• So far we have that

$$\textcircled{1} \quad u = \sum_i c_i \phi_i \quad \text{with} \quad \left. \begin{array}{l} \Delta \phi_i = -\lambda_i \phi_i \quad \text{in } \Omega \\ \phi_i = 0 \quad \text{on } \partial\Omega \end{array} \right\} \textcircled{2}$$

We can use orthogonality as usual to find c_i in terms of u (but now this is our unknown) and ϕ_i :

$$u = \sum c_i \phi_i \Rightarrow c_i = \frac{\int_{\Omega} u \phi_i}{\int_{\Omega} (\phi_i)^2}.$$

We use now ② to substitute on the numerator

$$\phi_i = \frac{-1}{\lambda_i} \Delta \phi_i, \text{ so that}$$

$$c_i = \frac{-1}{\lambda_i \int_{\Omega} (\phi_i)^2} \int_{\Omega} u \Delta \phi_i,$$

and now we integrate by parts twice,

$$c_i = \frac{-1}{\lambda_i \int_{\Omega} (\phi_i)^2} \left(\underbrace{\int_{\Omega} \phi_i \Delta u}_{\substack{= \\ Q}} - \underbrace{\int_{\partial \Omega} \phi_i \vec{n} \cdot \nabla u}_{=0} + \underbrace{\int_{\partial \Omega} u \vec{n} \cdot \nabla \phi_i}_{f} \right),$$

so in conclusion

$$c_i = -\frac{1}{\lambda_i} \frac{\int_{\Omega} \phi_i Q + \int_{\partial \Omega} f \vec{n} \cdot \nabla \phi_i}{\int_{\Omega} (\phi_i)^2},$$

← this is the boundary term

we were missing
on page -224-.

$$u = \sum_i c_i \phi_i$$

the formula will be
different for different BC

Summary: $\Delta u = Q$ in Ω $\left\{ \begin{array}{l} \text{BCs} \dots \partial\Omega \end{array} \right.$
 (same type)

1) $u = \sum_i c_i \phi_i$ with $\left\{ \begin{array}{l} \Delta \phi_i = -\lambda_i \phi_i \text{ in } \Omega \\ \text{homogeneous BCs} \dots \partial\Omega \end{array} \right.$

2) $c_i = \frac{\int_{\Omega} u \phi_i}{\int_{\Omega} (\phi_i)^2} = \frac{-1}{\lambda_i} \frac{\int_{\Omega} u \Delta \phi_i}{\int_{\Omega} (\phi_i)^2} = \dots$ integrate by parts, use BCs, and substitute $\Delta u = Q$. //

Remark: A second option to solve these problems is to split them into two. For example,

$$\left. \begin{array}{l} \Delta u = Q \text{ in } \Omega \\ u = f \text{ on } \partial\Omega \end{array} \right\} \rightarrow \left\{ \begin{array}{l} \Delta u_1 = 0 \text{ in } \Omega \\ u_1 = f \text{ on } \partial\Omega \end{array} \right. \left\{ \begin{array}{l} \Delta u_2 = Q \text{ in } \Omega \\ u_2 = 0 \text{ on } \partial\Omega \end{array} \right.$$

with $u = u_1 + u_2$.

• Notice that from previous lectures we know how to solve for u_1 .

• For u_2 , because the BCs are already homogeneous, we can differentiate term by term in $u = \sum_i c_i \phi_i$! //